

# **A New 4-D Hyperchaotic System with a Two-Scroll Attractor, its Adaptive Control and Circuit Simulation**

**Sundarapandian Vaidyanathan**

Research and Development Centre

Vel Tech University

Chennai, Tamil Nadu, India

[sundarvtu@gmail.com](mailto:sundarvtu@gmail.com)

**Aceng Sambas**

Department of Mechanical Engineering

Universitas Muhammadiyah Tasikmalaya

Tasikmalaya, Indonesia

[acengs@umtas.ac.id](mailto:acengs@umtas.ac.id)

**Sukono**

Department of Mathematics,

Faculty of Mathematics and Natural Sciences

Universitas Padjadjaran, Indonesia.

[sukono@unpad.ac.id](mailto:sukono@unpad.ac.id)

**Subiyanto**

Department of Marine Science

Universitas Padjadjaran, Indonesia.

[subiyanto@unpad.ac.id](mailto:subiyanto@unpad.ac.id)

**Mustafa Mamat**

Faculty of Informatics and Computing

Universiti Sultan Zainal Abidin

Kuala Terengganu, Malaysia

[must@unisza.edu.my](mailto:must@unisza.edu.my)

**Abdul Talib Bon**

Department of Production and Operations

University Tun Hussein Onn Malaysia, Malaysia

[talibon@gmail.com](mailto:talibon@gmail.com)

## **Abstract**

In the chaos literature, there is great interest in the modelling of hyperchaotic systems. This paper announces a new 4-D hyperchaotic system with a two-scroll attractor. The phase portraits and qualitative properties of the new hyperchaotic system are explored with MATLAB plots, bifurcation diagram, Lyapunov exponents, Kaplan-Yorke dimension, etc. As a control application, the adaptive stabilization of the new hyperchaotic system is detailed using Lyapunov stability theory. Finally, an electronic circuit of the new hyperchaotic system with MultiSIM is designed and a good match between the plots of the theoretical model and the circuit model of the new hyperchaotic system is shown.

## Keywords

Hyperchaos, hyperchaotic systems, adaptive control, stabilization and electronic circuit.

## 1. Introduction

Chaos theory deals with chaotic systems which are nonlinear dynamical systems highly sensitive to small changes in the initial conditions (Azar and Vaidyanathan, 2016; Vaidyanathan and Volos, 2016). Chaos theory has several applications in science and engineering such as memristors (Vaidyanathan and Volos, 2017), weather systems (Lorenz, 1963; Vaidyanathan *et al.* 2015), biological systems (Vaidyanathan 2015a, 2015b, 2015c, 2015d, Voorsluijs and Decker, 2016), neural systems (Vaidyanathan 2015e, 2015f, 2015g, 2015h), oscillators (Pakiriswamy and Vaidyanathan, 2012; Rasappan and Vaidyanathan, 2012; Vaidyanathan 2012a, 2012b, 2014, 2015i, 2015j; Vaidyanathan and Pakiriswamy, 2011, 2012; Vaidyanathan and Rasappan, 2011; Vaidyanathan and Sampath, 2012), chemical reactors (Vaidyanathan, 2015k), circuits (Li *et al.*, 2014, 2015; Daltzis *et al.*, 2018; Mamat *et al.*, 2018; Pham *et al.*, 2016; Sambas *et al.*, 2018a, 2018b; Vaidyanathan, 2018a, 2018b, 2018c, 2018d, 2018e), finance (Tacha *et al.*, 2016), robotics (Sambas *et al.*, 2016; Vaidyanathan *et al.* 2017), etc.

Hyperchaotic systems are chaotic systems having two or more positive Lyapunov exponents, which means that the state evolution for hyperchaotic systems can expand in two or more directions yielding high complexity (Azar and Vaidyanathan, 2016; Vaidyanathan and Volos, 2016). Hyperchaotic systems are suitable for many engineering applications such as steganography (Vaidyanathan *et al.*, 2018f), cryptography (Murali, 2000; Wu *et al.*, 2014), image encryption (Xue *et al.*, 2018), secure communications (Li *et al.*, 2005), etc.

This work announces a new 4-D hyperchaotic system with a two-scroll attractor. It is interesting to know that the new hyperchaotic system has only two quadratic nonlinearities and it has two positive Lyapunov exponents confirming its hyperchaotic behavior. We show that the new hyperchaotic system has an unstable, saddle-point, equilibrium at the origin. We describe the dynamic properties of the new hyperchaotic system via phase plots, Lyapunov exponents, Kaplan-Yorke dimension, etc. As a control application, we build an adaptive feedback control law for globally stabilizing the trajectories of the new hyperchaotic system. Finally, an electronic circuit of the new hyperchaotic system with two-scroll attractor is designed via MultiSIM to verify the feasibility of the theoretical model. Our simulations show a good match between the MATLAB plots and the MultiSIM simulations of the new hyperchaotic system with a two-scroll attractor.

## 2. A New 4-D Hyperchaotic System with a Two-Scroll Attractor

In this work, we report a new 4-D system given by the dynamics

$$\begin{cases} \dot{x}_1 = a(x_2 - x_1) + x_4 \\ \dot{x}_2 = bx_1 - x_2 - dx_1x_3 + x_4 \\ \dot{x}_3 = x_1x_2 - cx_3 \\ \dot{x}_4 = -x_1 - x_2 \end{cases} \quad (1)$$

In (1),  $X = (x_1, x_2, x_3, x_4)$  is the state and  $a, b, c, d$  are constant parameters. We shall show that the system (1) exhibits a hyperchaotic two-scroll attractor for the parameter values

$$a = 12, \quad b = 40, \quad c = 5, \quad d = 12 \quad (2)$$

The Lyapunov exponents of the new system (1) for  $(a, b, c, d) = (12, 40, 5, 12)$  and  $X(0) = (0.2, 0.2, 0.2, 0.2)$  are calculated for  $T = 10^5$  seconds using Wolf's algorithm (Wolf *et al.*, 1985) as follows:

$$LE_1 = 1.0905, \quad LE_2 = 0.1278, \quad LE_3 = 0, \quad LE_4 = -19.1710 \quad (3)$$

From (3), we note that the new 4-D system (1) is hyperchaotic since it has two positive Lyapunov exponents. Also, the new system (1) is dissipative since the sum of the Lyapunov exponents is negative.

Furthermore, the Kaplan-Yorke dimension of the new system (1) is determined as follows:

$$D_{KY} = 3 + \frac{LE_1 + LE_2 + LE_3}{|LE_4|} = 3.0635 \quad (4)$$

The value of the Kaplan-Yorke dimension obtained in (4) shows the high complexity of the new hyperchaotic system (4), which is dissipative with a strange attractor.

We note that the new hyperchaotic system (1) remains invariant under the change of coordinates given by

$$(x_1, x_2, x_3, x_4) \mapsto (-x_1, -x_2, x_3, -x_4) \quad (5)$$

This shows that the new hyperchaotic system (1) has rotation symmetry about the  $x_3$  – coordinate axis. Hence, all the non-trivial trajectories of the new hyperchaotic system (1) must have twin trajectories associated with them.

We get the equilibrium points of the new hyperchaotic system (1) by solving the following equations for the hyperchaotic case  $(a, b, c, d) = (12, 40, 5, 12)$ :

$$a(x_2 - x_1) + x_4 = 0 \quad (6a)$$

$$bx_1 - x_2 - dx_1x_3 + x_4 = 0 \quad (6b)$$

$$x_1x_2 - cx_3 = 0 \quad (6c)$$

$$-x_1 - x_2 = 0 \quad (6d)$$

Solving the system (6), it is easy to see that the unique equilibrium of the new hyperchaotic system (1) is given by  $E_0 = (0, 0, 0, 0)$ .

The Jacobian of the new hyperchaotic system (1) at  $E_0$  for  $(a, b, c, d) = (12, 40, 5, 12)$  is obtained as follows.

$$J = \begin{bmatrix} -12 & 12 & 0 & 1 \\ -40 & -1 & 0 & 1 \\ 0 & 0 & -5 & 0 \\ -1 & -1 & 0 & 0 \end{bmatrix} \quad (7)$$

The eigenvalues of  $J$  are obtained as  $\lambda_1 = 0.0303$ ,  $\lambda_2 = -5$ ,  $\lambda_{3,4} = -6.5152 \pm 21.2591i$ .

This shows that the equilibrium point  $E_0$  is a saddle-point, which is unstable.

Hence, the strange attractor associated with the new hyperchaotic system (1) is a self-excited attractor.

The phase plots of the new hyperchaotic system (1) for  $X(0) = (0.2, 0.2, 0.2, 0.2)$  and  $(a, b, c, d) = (12, 40, 5, 12)$  are shown in Figures 1-4. From the phase plots, the new hyperchaotic system (1) has a two-scroll hyperchaotic attractor. Also, Figure 5 shows the Lyapunov exponents of the new hyperchaotic system (1). The presence of two positive Lyapunov exponents in Figure 5 confirms that the new 4-D system (1) is indeed a hyperchaotic system.

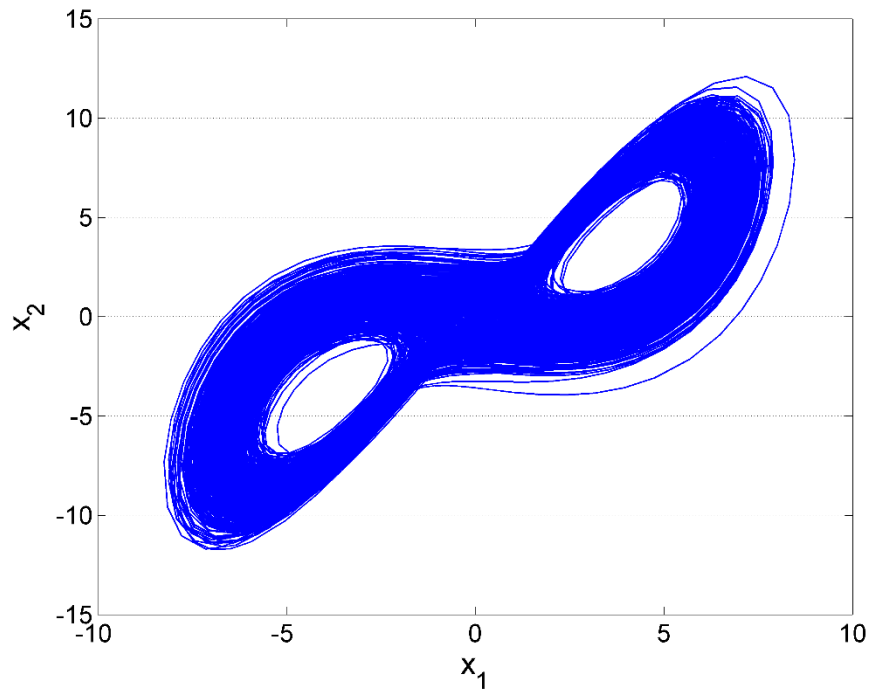


Figure 1. 2-D phase plot of the new hyperchaotic system (1) in the  $(x_1, x_2)$  - plane for  $X(0) = (0.2, 0.2, 0.2, 0.2)$  and  $(a, b, c, d) = (12, 40, 5, 12)$

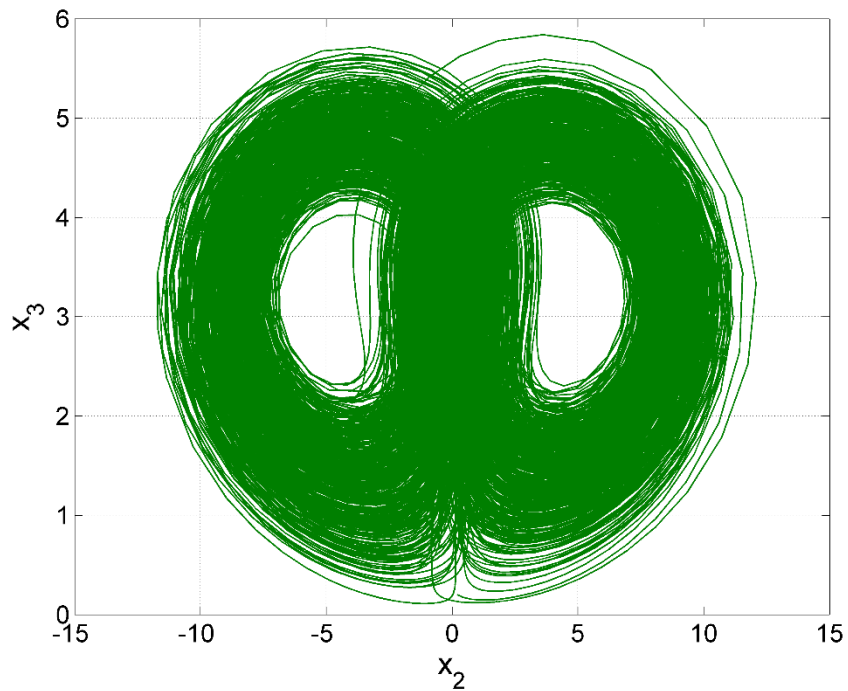


Figure 2. 2-D phase plot of the new hyperchaotic system (1) in the  $(x_2, x_3)$  - plane for  $X(0) = (0.2, 0.2, 0.2, 0.2)$  and  $(a, b, c, d) = (12, 40, 5, 12)$

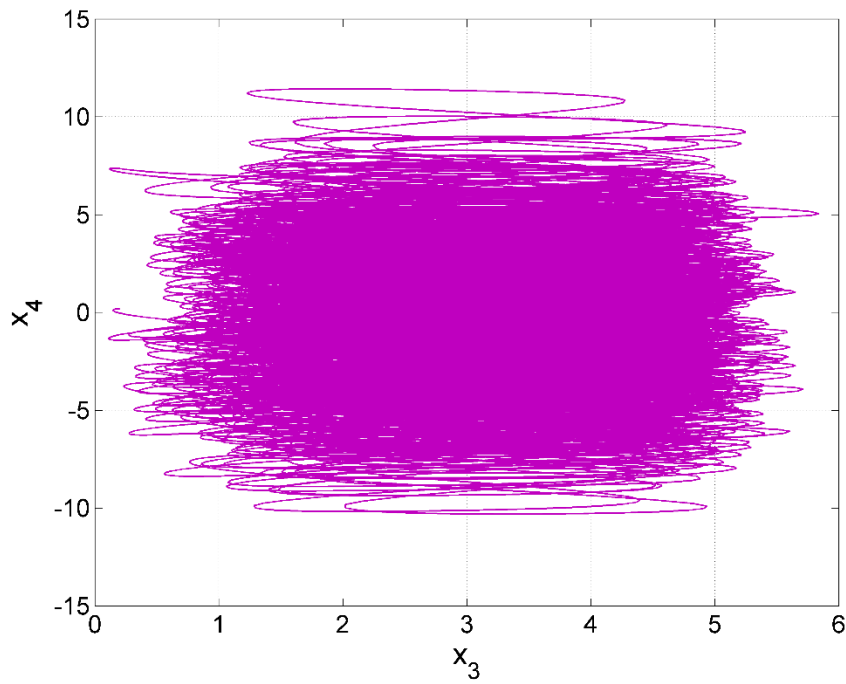


Figure 3. 2-D phase plot of the new hyperchaotic system (1) in the  $(x_3, x_4)$  – plane for  $X(0) = (0.2, 0.2, 0.2, 0.2)$  and  $(a, b, c, d) = (12, 40, 5, 12)$

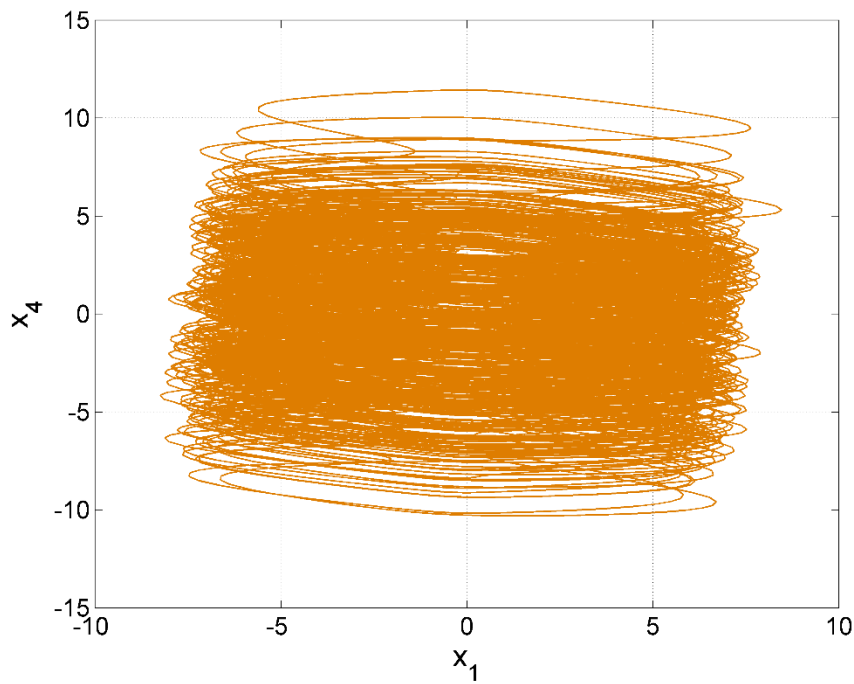


Figure 4. 2-D phase plot of the new hyperchaotic system (1) in the  $(x_1, x_4)$  – plane for  $X(0) = (0.2, 0.2, 0.2, 0.2)$  and  $(a, b, c, d) = (12, 40, 5, 12)$

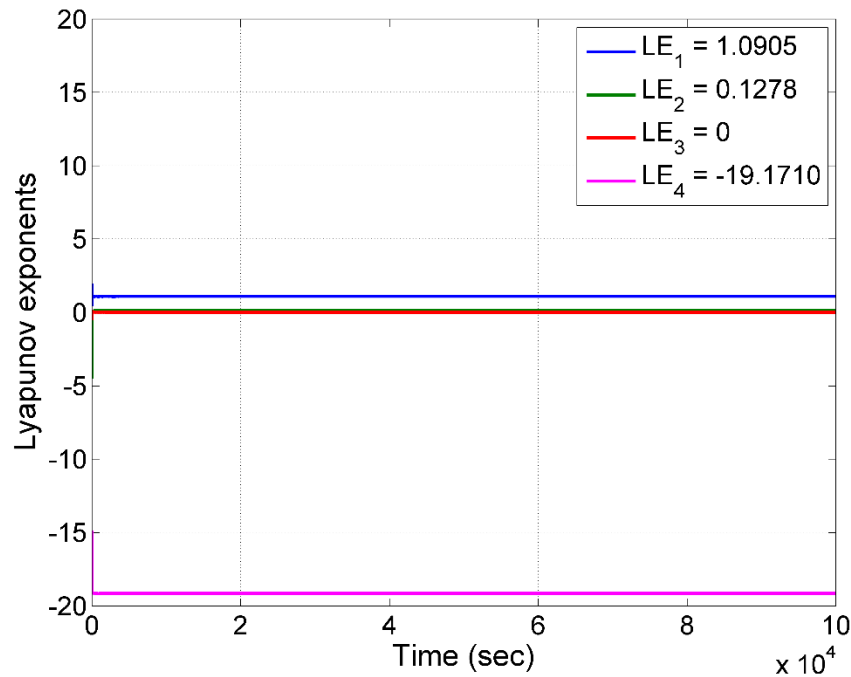


Figure 5. Lyapunov exponents of the new system (1) for  $X(0) = (0.2, 0.2, 0.2, 0.2)$  and  $(a, b, c, d) = (12, 40, 5, 12)$

### 3. Global Chaos Control of the New 4-D Hyperchaotic System with a Two-Scroll Attractor

In this section, we consider the problem of globally stabilizing the trajectories of the new 4-D hyperchaotic system, which is given by the controlled dynamics

$$\begin{cases} \dot{x}_1 = a(x_2 - x_1) + x_4 + u_1 \\ \dot{x}_2 = bx_1 - x_2 - dx_1x_3 + x_4 + u_2 \\ \dot{x}_3 = x_1x_2 - cx_3 + u_3 \\ \dot{x}_4 = -x_1 - x_2 + u_4 \end{cases} \quad (8)$$

In (8),  $X = (x_1, x_2, x_3, x_4)$  is the state,  $u = (u_1, u_2, u_3, u_4)$  is the adaptive control to be designed and  $a, b, c, d$  are system parameters which are not available for measurement. So, we use estimates  $A(t), B(t), C(t), D(t)$  in lieu of the unknown system parameters  $a, b, c, d$ , respectively.

We consider the adaptive control given by

$$\begin{cases} u_1 = -A(t)(x_2 - x_1) - x_4 - k_1x_1 \\ u_2 = -B(t)x_1 + x_2 + D(t)x_1x_3 - x_4 - k_2x_2 \\ u_3 = -x_1x_2 + C(t)x_3 - k_3x_3 \\ u_4 = x_1 + x_2 - k_4x_4 \end{cases} \quad (9)$$

In (9),  $k_1, k_2, k_3, k_4$  are positive constants, and  $A(t), B(t), C(t), D(t)$  are parameter estimates. We define the parameter estimation error as follows:

$$\begin{cases} e_a = a - A(t) \\ e_b = b - B(t) \\ e_c = c - C(t) \end{cases} \quad (10)$$

Differentiating (10), we obtain the following:

$$\begin{cases} \dot{e}_a = -\dot{A}(t) \\ \dot{e}_b = -\dot{B}(t) \\ \dot{e}_c = -\dot{C}(t) \end{cases} \quad (11)$$

By substituting the adaptive control law (9) into (8), we get the closed-loop state dynamics as follows:

$$\begin{cases} \dot{x}_1 = [a - A(t)](x_2 - x_1) - k_1 x_1 \\ \dot{x}_2 = [b - B(t)]x_1 - [d - D(t)]x_1 x_3 - k_2 x_2 \\ \dot{x}_3 = -[c - C(t)]x_3 - k_3 x_3 \\ \dot{x}_4 = -k_4 x_4 \end{cases} \quad (12)$$

Using the definition (10), it is possible to simplify the closed-loop dynamics (12) as follows:

$$\begin{cases} \dot{x}_1 = e_a(x_2 - x_1) - k_1 x_1 \\ \dot{x}_2 = e_b x_1 - e_d x_1 x_3 - k_2 x_2 \\ \dot{x}_3 = -e_c x_3 - k_3 x_3 \\ \dot{x}_4 = -k_4 x_4 \end{cases} \quad (13)$$

We consider the quadratic Lyapunov function  $V$  defined as follows:

$$V(X, e_a, e_b, e_c, e_d) = \frac{1}{2}(x_1^2 + x_2^2 + x_3^2 + x_4^2) + \frac{1}{2}(e_a^2 + e_b^2 + e_c^2 + e_d^2) \quad (14)$$

Differentiating  $V$  along the trajectories of (13) and (11), we obtain the following:

$$\dot{V} = -\sum_{i=1}^4 k_i e_i^2 + e_a [x_1(x_2 - x_1) - \dot{A}] + e_b [x_1 x_2 - \dot{B}] + e_c [-x_3^2 - \dot{C}] + e_d [-x_1 x_2 x_3 - \dot{D}] \quad (15)$$

In view of (15), we propose the following parameter update law:

$$\begin{cases} \dot{A} = x_1(x_2 - x_1) \\ \dot{B} = x_1 x_2 \\ \dot{C} = -x_3^2 \\ \dot{D} = -x_1 x_2 x_3 \end{cases} \quad (16)$$

Thus, we prove the following main result of this section.

The new result is established using Lyapunov stability theory (Khalil, 2002).

**Theorem 1.** The new hyperchaotic two-scroll system (8) with unknown system parameters  $a, b, c, d$  is globally stabilized for all initial conditions  $X(0) \in R^4$  with the adaptive feedback control law (9) and the parameter update law (16), where the control gains  $k_1, k_2, k_3, k_4$  are positive constants.

**Proof.** First, we note that the candidate Lyapunov function  $V$  given in Eq. (14) is a quadratic and positive definite function defined on  $R^8$ . Next, we substitute the parameter update law (16) into Eq. (15). This yields the following:

$$\dot{V} = -k_1 e_1^2 - k_2 e_2^2 - k_3 e_3^2 - k_4 e_4^2 \quad (17)$$

This shows that  $\dot{V}$  is a negative semi-definite function on  $R^8$ . Hence, by Barbalat's Lemma (Khalil, 2002), we conclude that the controlled state  $x_i(t), (i = 1, 2, 3, 4)$  of the closed-loop system (13) exponentially converges to zero as  $t \rightarrow \infty$ . This completes the proof. ■

For numerical simulations, we take the control gains as  $k_i = 8$  for  $i = 1, 2, 3, 4$ . Also, we take the parameter values as in the hyperchaotic case (2), i.e.  $(a, b, c, d) = (12, 40, 5, 12)$ .

As the initial state, we take  $X(0) = (1.5, 8.3, 4.1, 9.7)$ . As the initial state for the parameter estimate, we take  $(A(0), B(0), C(0), D(0)) = (8, 4, 5, 2)$ .

Figure 6 shows the time-history of the state  $X(t)$  of the controlled hyperchaotic two-scroll system (8), when the adaptive feedback control law (9) and the parameter update law (16) are implemented.

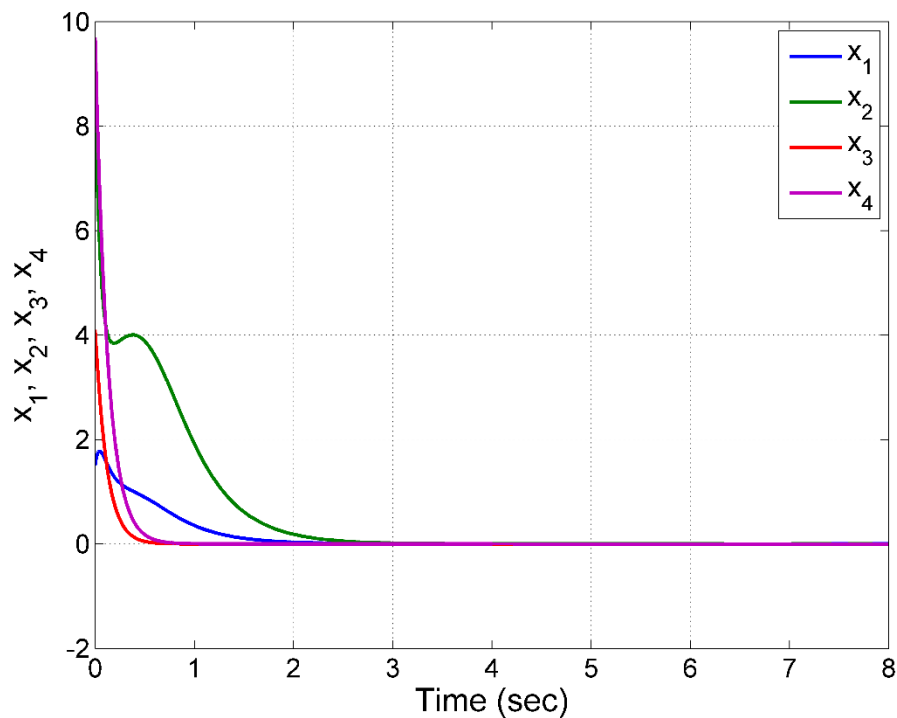


Figure 6. Time-history of the state  $X(t)$  of the controlled system (8)



#### 4. Circuit Simulation of the New 4-D Hyperchaotic System with a Two-Scroll Attractor

In this section, electronic circuit of the hyperchaotic system (1) is implemented. Figure 7 presents the schematic diagram of the proposed analog simulator for hyperchaotic system (1). According to Kirchhoff's circuit laws, the equations in terms of the circuit parameters are:

$$\begin{cases} \dot{x}_1 = \frac{1}{C_1 R_1} x_2 - \frac{1}{C_1 R_2} x_1 + \frac{1}{C_1 R_3} x_4 \\ \dot{x}_2 = \frac{1}{C_2 R_4} x_1 - \frac{1}{C_2 R_5} x_2 - \frac{1}{10 C_2 R_6} x_1 x_3 + \frac{1}{C_2 R_7} x_4 \\ \dot{x}_3 = \frac{1}{10 C_3 R_8} x_1 x_2 - \frac{1}{C_3 R_9} x_3 \\ \dot{x}_4 = -\frac{1}{C_4 R_{10}} x_1 - \frac{1}{C_4 R_{11}} x_2 \end{cases} \quad (18)$$

In (18),  $x_1$ ,  $x_2$ ,  $x_3$  and  $x_4$  correspond to the voltages on the integrators U1A, U2A, U3A and U4A, respectively. The values of the circuit components are chosen as:  $R_1 = R_2 = 33.33 \text{ k}\Omega$ ,  $R_3 = R_5 = R_7 = R_{10} = R_{11} = 400 \text{ k}\Omega$ ,  $R_4 = 10 \text{ k}\Omega$ ,  $R_6 = 3.33 \text{ k}\Omega$ ,  $R_8 = 40 \text{ k}\Omega$ ,  $R_9 = 80 \text{ k}\Omega$ ,  $R_{12} = R_{13} = R_{14} = R_{15} = R_{16} = R_{17} = 100 \text{ k}\Omega$ ,  $C_1 = C_2 = C_3 = C_4 = 1 \text{ nF}$ . Phase portraits from the output of the integration channels are given in Figure 8. We notice a very good similarity between numerical simulations (see Figures 1-4) and oscilloscope results (see Figure 8).

#### 5. Conclusions

In this paper, we proposed a new 4-D hyperchaotic system with a two-scroll attractor. Lyapunov exponents are derived to prove that the system is hyperchaotic. Also, the numerical simulation and electronic circuit design of a new 4-D hyperchaotic system with a two-scroll attractor are fulfilled. As a control application, we devised adaptive control law for globally stabilizing all the trajectories of the new hyperchaotic system with a two-scroll attractor. Hyperchaotic systems have several applications in engineering areas such as cryptography, encryption, secure communications, etc. It is hoped that the new hyperchaotic system is applied to many engineering areas.

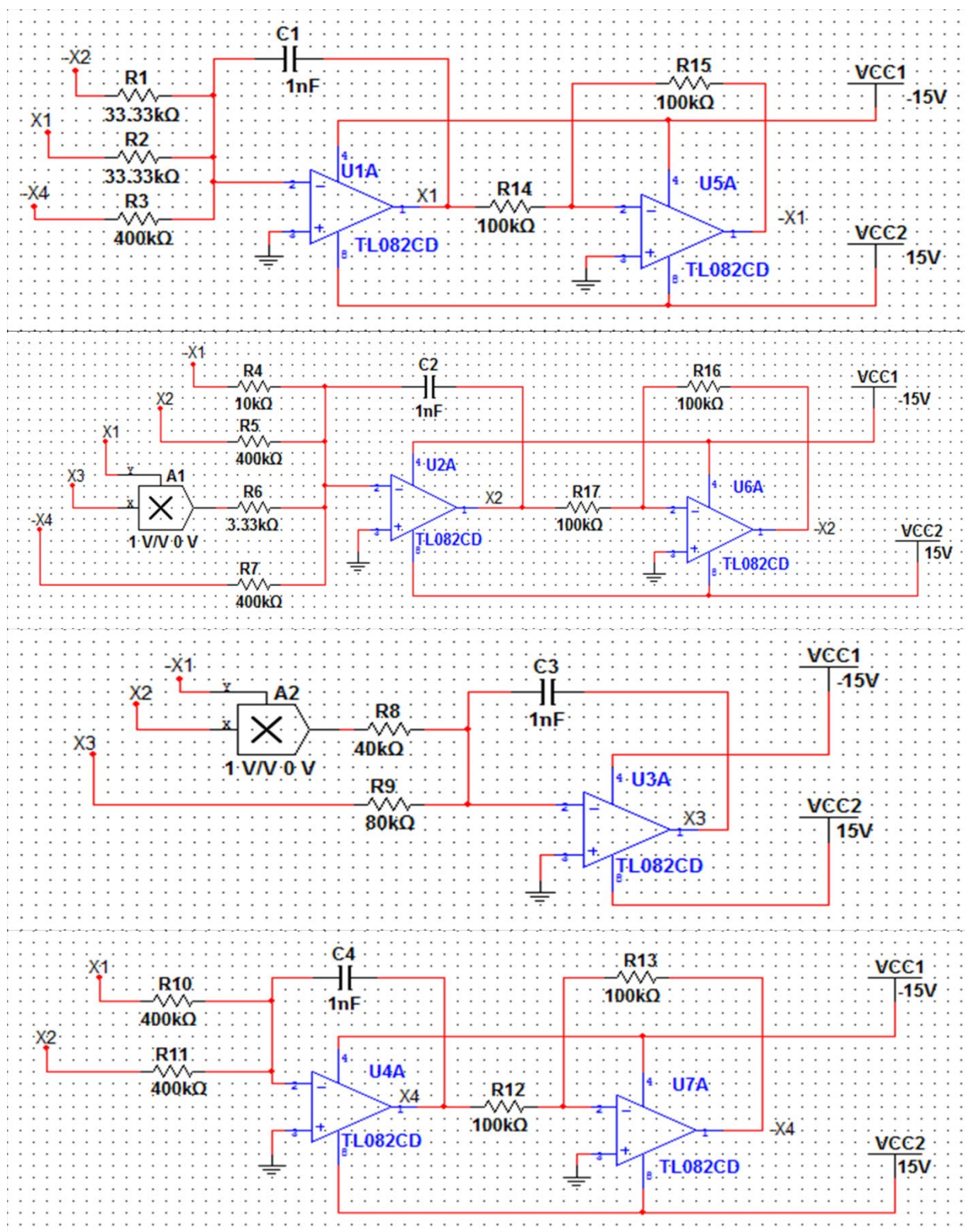
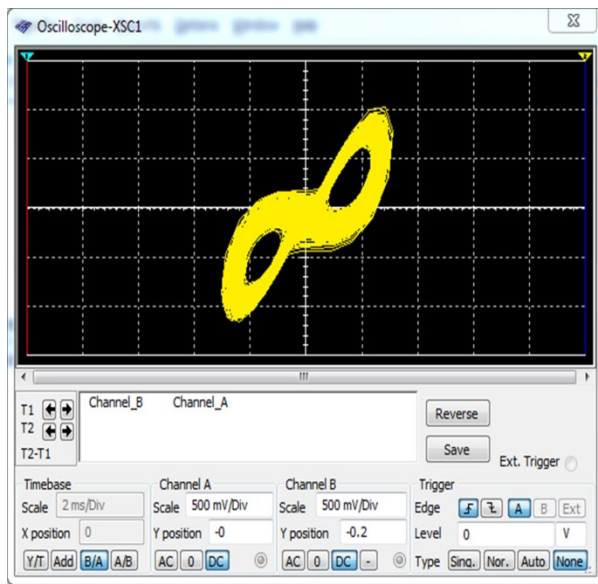
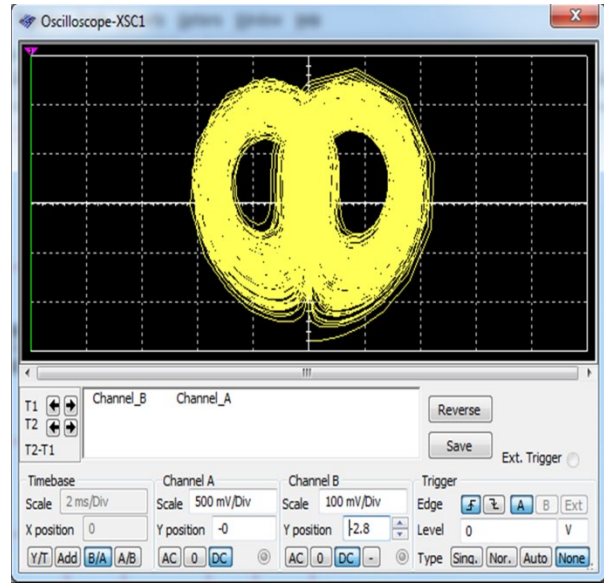


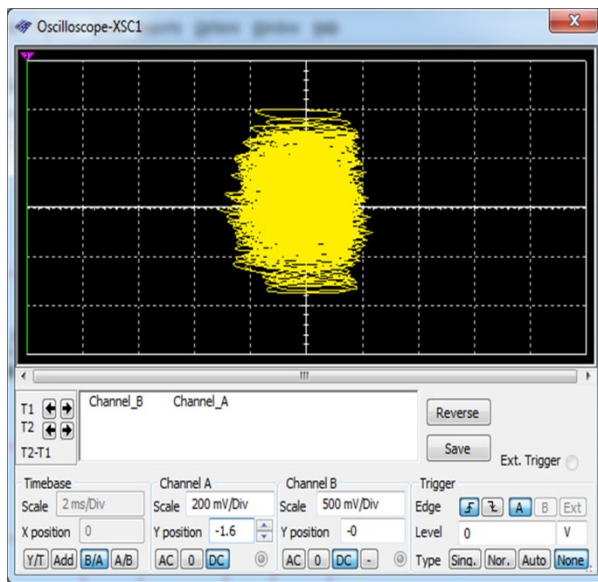
Figure 7. The electronic circuit schematic of the new 4-D hyperchaotic system with a two-scroll attractor



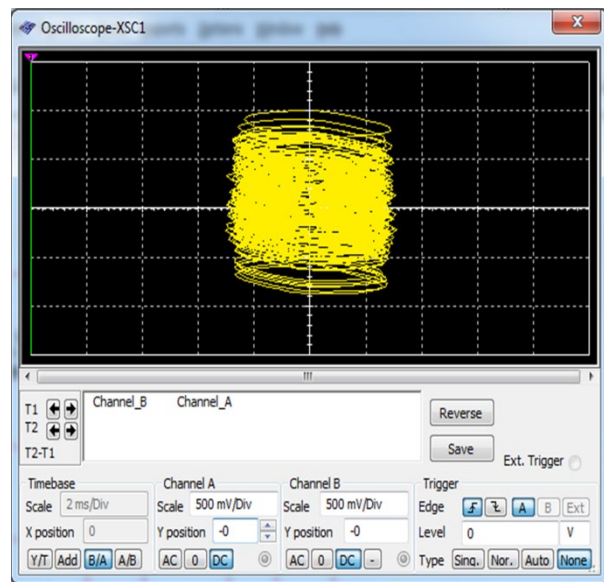
(a)



(b)



(c)



(d)

Figure 8 MultiSIM chaotic attractors of the new 4-D hyperchaotic system with a two-scroll attractor  
(a)  $x_1$ -  $x_2$  plane, (b)  $x_2$ -  $x_3$  plane, (c)  $x_3$ -  $x_4$  plane and (d)  $x_1$ -  $x_4$  plane.

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## **Biographies**

**Sundarapandian Vaidyanathan** is a Professor and Dean at the Research and Development Centre, Vel Tech University, Chennai, India. He received his D.Sc in Electrical and Systems Engineering from Washington University, St. Louis, USA in 1996. He has published over 460 Scopus-indexed research papers. His current research focuses on control systems, chaos theory, chaotic and hyperchaotic systems, sliding mode control, neuro-fuzzy control, computational science, circuits and memristors. He is the Editor-in-Chief of International Journal of Nonlinear Dynamics and Control (IJNDC), Inderscience Publishers, Olney, UK. He is also in the Editorial Boards of many control journals published by Inderscience, Olney, UK.

**Aceng Sambas** is currently a Lecturer at the Muhammadiyah University of Tasikmalaya, Indonesia since 2015. He received his MSc in Mathematics from the Universiti Sultan Zainal Abidin (UniSZA), Malaysia in 2015. His current research focuses on dynamical systems, chaotic signals, electrical engineering, computational science, signal processing, robotics, embedded systems and artificial intelligence.

**Sukono** is a lecturer in the Department of Mathematics, Faculty of Mathematics and Natural Sciences, Universitas Padjadjaran. Currently serves as Head of Master's Program in Mathematics, the field of applied mathematics, with a field of concentration of financial mathematics and actuarial sciences.

**Subiyanto** is a lecturer in the Department of Marine Science, Faculty of Fishery and Marine Science, Universitas Padjadjaran. He received his Ph.D in School of Ocean Engineering from Universiti Malaysia Terengganu (UMT), Malaysia in 2017. His research focuses on applied mathematics, numerical analysis and computational science.

**Mustafa Mamat** is currently a Professor and the Dean of Graduate School at Universiti Sultan Zainal Abidin (UniSZA), Malaysia since 2013. He was first appointed as a Lecturer at the Universiti Malaysia Terengganu (UMT) in 1999. He obtained his PhD from the UMT in 2007 with specialization in optimization. Later on, he was appointed as a Senior Lecturer in 2008 and then as an Associate Professor in 2010 also at the UMT. To date, he has successfully supervised more than 60 postgraduate students and published more than 200 research papers in various international journals and conferences. His research interests include conjugate gradient methods, steepest descent methods, Broyden's family and quasi-Newton methods.

**Abdul Talib Bon** is a Professor of Production and Operations Management in the Faculty of Technology Management and Business at the Universiti Tun Hussein Onn Malaysia since 1999. He has a PhD in Computer Science, which he obtained from the Universite de La Rochelle, France in the year 2008. His doctoral thesis was on topic Process Quality Improvement on Beltline Moulding Manufacturing. He studied Business Administration in the Universiti Kebangsaan Malaysia for which he was awarded the MBA in the year 1998. He's bachelor degree and diploma in Mechanical Engineering which he obtained from the Universiti Teknologi Malaysia. He received his postgraduate certificate in Mechatronics and Robotics from Carlisle, United Kingdom in 1997. He had published more 150 International Proceedings and International Journals and 8 books. He is a member of MSORSM, IIF, IEOM, IIE, INFORMS, TAM and MIM.